

Thermo Mechanical Model for Orthogonal Metal Cutting

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ABSTRACT: A finite element model of the orthogonal metal cutting process with a coupled thermal-elastic-plastic constitutive equation under large deformations based on update Lagrangian formulation is presented. Plane strain conditions with a flow stress function of strain, strain rate, work hardening and temperature is assumed. Principle of virtual work is applied in order to construct an equilibrium rate solution. The problem is formulated, discretized and the resulting equations solved with the finite element method. As an application the temperature field is computed for the case of machining AISI 1020 steel. The stress field is also shown in order to present the coupled problem. The interface problem is modeled with a simple Coulomb friction surface. The model produced good agreement with results presented in the literature.

1 INTRODUCTION

Metal cutting is a quite complex phenomenon, of multi physics type, where elasticity and plasticity, fracture, contacts, heat transfer, among others takes place simultaneously. Much work has been devoted to its understanding. Initially having an experimental character and afterwards theoretical as well. Since the forties several attempts were made pursuing a linking between empirical knowledge and theoretical modeling. Such attempts begun at forties (Piispanen, 1948) and (Merchant 1945), concentrating in explain the mechanisms of the chip formation. A recent review can be seen in (Mackerle 1999).

In metal cutting the material is subjected to extremely high strains in two principal regions: the shear zone or primary deformation zone, and another zone, known as secondary deformation zone. Determination of stresses and temperatures in this region is a central problem in understanding machining processes, as the problem is of a coupled one.

Many models were developed to address this solution. Of particular success were the ones considering the finite element tool, using update Lagrange scheme. Here it is presented some details of one such a solution, regarding some previous work in the area (Lin, 1990, 1992, 1995), (Komvopoulos, 1991), (Tay 1976 and Stevenson 1983), (Rebello & Kobayashi, 1980) and (Muraka & all, 1979). A quasi-static solution with typical machining parameters will be formulated and implemented in this work.

2 FORMULATION

2.1 Orthogonal Cutting Model

The chip formation model to analyze has the geometry shown in Figure 1. In it is possible to distinguish two regions separated by a division line AB, the cutting line, marking the positions separation is going to take place. Dimensions for this part of the workpiece are shown in mm, in the same figure. It corresponds to a region of the workpiece where deformation is localized. For the mechanical and thermal problems, boundary conditions shown in Figure 2 apply. In the left lower side, it was prescribed null velocities, for the horizontal, 1, and vertical, 2, directions. Temperatures have a fixed value T_0 .

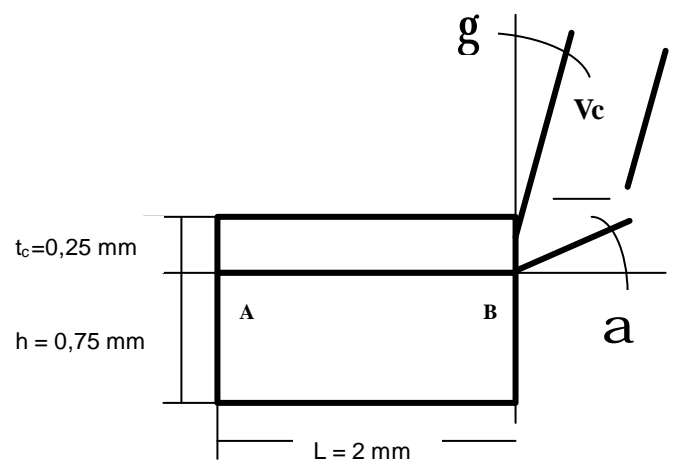


Figure 1. Chip formation model

The virtual velocity field is assumed to be independent and obeying the boundary conditions. The nominal, or Lagrange stress is denoted by \mathbf{s} , the nominal traction on S_o by \mathbf{f}_o and body force by on V_o by \mathbf{b}_o (Yamada, 1973). If we consider the expressions above and perform the indicated operations we may write that

$$\delta \dot{W} = \int_{V_o} \mathbf{F}^{-1} [\mathbf{s} + \mathbf{W} \cdot \mathbf{s} - \mathbf{s} \cdot \mathbf{W} - \mathbf{L}_m \mathbf{s} + \text{tr} \mathbf{L}_m \mathbf{s}] \quad (10)$$

$$\mathbf{W} = \frac{1}{2}(\mathbf{L} + \mathbf{L}^T) \quad \mathbf{D} = \frac{1}{2}(\mathbf{L} + \mathbf{L}^T) \quad \mathbf{L} = \frac{\partial \mathbf{v}}{\partial \mathbf{x}}; \mathbf{v} = \delta \dot{\mathbf{x}}$$

Where its geometric part may have the velocity gradient \mathbf{L} factored out through a matrix \mathbf{D}^g (Obikawa, 1966). Discretized the part with finite elements, being the nodal velocities \mathbf{V} the independent variables, results in the matrix form:

$$(\mathbf{K}^{ep} + \mathbf{K}^g) \mathbf{V} = \mathbf{F} \delta \dot{\mathbf{x}} \\ \mathbf{K}^{ep} = \int_{V_i} \mathbf{B}_L^T \mathbf{C}^{ep} \mathbf{B}_D dV; \quad \mathbf{K}^g = \int_{V_i} \mathbf{B}_L^T \mathbf{D}^g \mathbf{B}_L dV \quad (11)$$

With the loading rate term composed of the parts:

$$\mathbf{f}_0^{\delta \dot{\mathbf{x}}} = \int_{S_o} \mathbf{N}^T \mathbf{f}_o^{\delta \dot{\mathbf{x}}} dS + \int_{V_o} \mathbf{N}^T \mathbf{b}_o^{\delta \dot{\mathbf{x}}} dV; \quad \mathbf{v} = \mathbf{N} \mathbf{V} \\ \mathbf{f}_0^{\delta \dot{\mathbf{x}}} = \int_{V_i} \mathbf{B}_L^T (\mathbf{C}^{ep} \mathbf{D}^i - \frac{\partial R}{\partial \mathbf{q}} \mathbf{q}^{\delta \dot{\mathbf{x}}}) dV; \quad \mathbf{D}^i = (\mathbf{a} \mathbf{q}) \mathbf{I} \quad (12) \\ \mathbf{f}_0^{\delta \dot{\mathbf{x}}} = - \int_{V_i} \mathbf{B}_L^T \frac{\partial R}{\partial \mathbf{e}} \mathbf{e}^{\delta \dot{\mathbf{x}}} dV; \quad \mathbf{e}^{\delta \dot{\mathbf{x}}} = \sqrt{\frac{2}{3} \mathbf{D} : \mathbf{D}}; \quad \mathbf{B}_L = \frac{\partial \mathbf{N}}{\partial \mathbf{x}}$$

Where α , thermal expansion and θ , temperature, couple the problem.

2.4 Unsteady State Thermal Conduction Analysis.

The integral form of the energy equation to be solved is given by the following expression

$$\int_{V_i} \mathbf{r} c \frac{\partial \mathbf{q}}{\partial t} dV = \int_{S_i} \mathbf{k} \frac{\partial \mathbf{q}}{\partial \mathbf{x}} \cdot \mathbf{n} dS + \mathcal{W}_f^{\delta \dot{\mathbf{x}}} + \mathcal{W}_p^{\delta \dot{\mathbf{x}}} \quad (13)$$

Where \mathbf{r} is the density, c is the specific heat, \mathbf{k} is the thermal conductivity, and W is the rate of dissipation due to friction and plasticity effects:

$$\mathcal{W}^{\delta \dot{\mathbf{x}}} = \int_{S_i} \mathbf{m} \mathbf{t} \cdot \delta \dot{\mathbf{x}} dS; \quad \mathcal{W}_p^{\delta \dot{\mathbf{x}}} = \int_{V_i} \mathbf{h} \bar{\mathbf{S}} \bar{\mathbf{D}}^p dV \quad (14)$$

Where η measures the part of the plastic work converted into heat while μ represents the friction coefficient associated with tangential forces \mathbf{t} . On applying the principle of virtual temperatures and upon discretizing the same volume with the above

interpolation function (Zienkiewicz and Taylor, 1991), as:

$$\mathbf{C} \mathbf{F} + \mathbf{K}_k \mathbf{T} = \mathcal{Q}_p^{\delta \dot{\mathbf{x}}} + \mathcal{Q}_f^{\delta \dot{\mathbf{x}}} + \mathcal{Q}_k^{\delta \dot{\mathbf{x}}} \\ \mathbf{C} = \int_{V_i} \mathbf{N}^T c \mathbf{N} dV; \quad \mathbf{K}_k = \int_{V_i} \mathbf{B}^T \mathbf{k} \mathbf{B} dV_i \quad (15) \\ \mathcal{Q}_p^{\delta \dot{\mathbf{x}}} = \int_{V_i} \mathbf{N}^T \mathbf{h} \mathcal{W}_p^{\delta \dot{\mathbf{x}}} dV; \quad \mathcal{Q}_f^{\delta \dot{\mathbf{x}}} = \int_{S_i^{\delta \dot{\mathbf{x}}}} \mathbf{m} \mathcal{W}_f^{\delta \dot{\mathbf{x}}} dS$$

Friction dissipation at the interface in slip condition.

2.5 Contact Problem

Solution of the contact problem, between tool and workpiece is based on Coulomb's law. In it, localized stick or slip conditions depend on the magnitude of the friction coefficient. For conditions of rigid stick, dissipation potential is added to the general virtual work statement used above. In the contact interface adhesion and slipping/sticking contact occur. Considered the discretization used above, the final result is the addition of stiffness term \mathbf{K}_f to the left side of expression and a vector term to the right side (Cheng and Kikuchi, 1985).

2.6 Separation Criterion

As the tool advances, each node on the path line was assumed to separate at the cutting edge into a chip surface node and a machined surface node when the strain energy density reached a critical value (Lin, 1992).

3. MODEL IMPLEMENTATION

A quasi-static simulation, including the model parameters considered above was implemented and results are shown ahead. The model was run assuming a HSS tool having a rake angle of 8 degrees, moving with a velocity of 122 m/min. Coefficient of friction between 0.2 and 0.5 were considered. The workpiece was assumed made of AISI 1020, $E=207$ GPa, $\nu=0.292$, $\sigma_0=210$ MPa. The model was implemented using an ABAQUS FEM code (ABAQUS, 1994). Chip formation and separation during the process is shown in Figure 3. Figure 4 shows the equivalent plastic strains. It concentrates in the primary shear zone and close to the tool-chip interface, in agreement with the Merchant's theory. Figure 6 shows the isotherms of the temperature distribution in the tool, chip and workpiece. It shows the temperature rise due to the combined effects of two principal heat sources, the plastic strain in the shear plane and the frictional heat dissipation in the chip-tool interface. In generating this plot it was assumed room temperature as initial condition, T_0 . Maximum temperatures occur inside the shear zone and at the secondary shear zone.

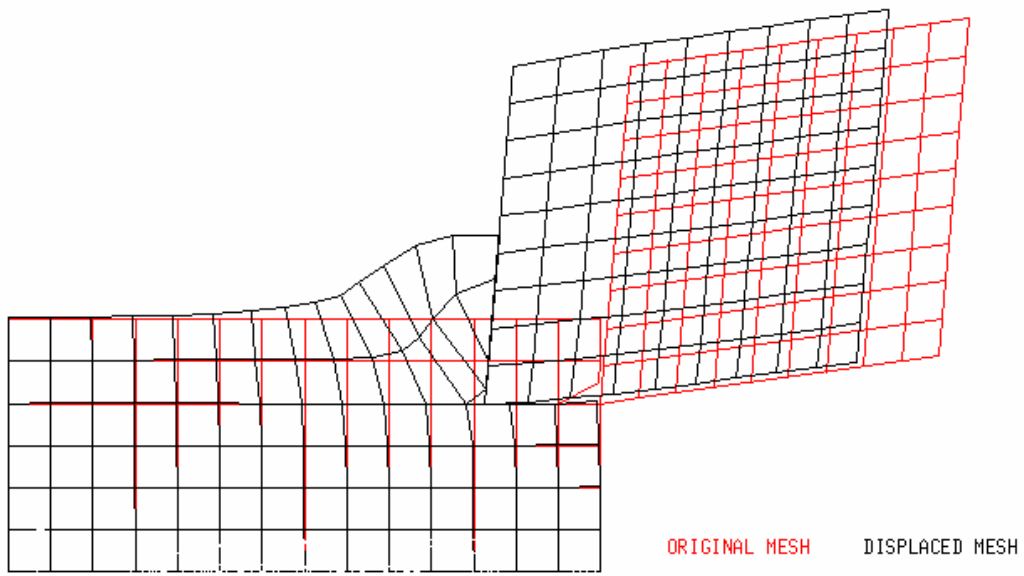


Figure 3. Chip formation: (red) undeformed, black, deformed

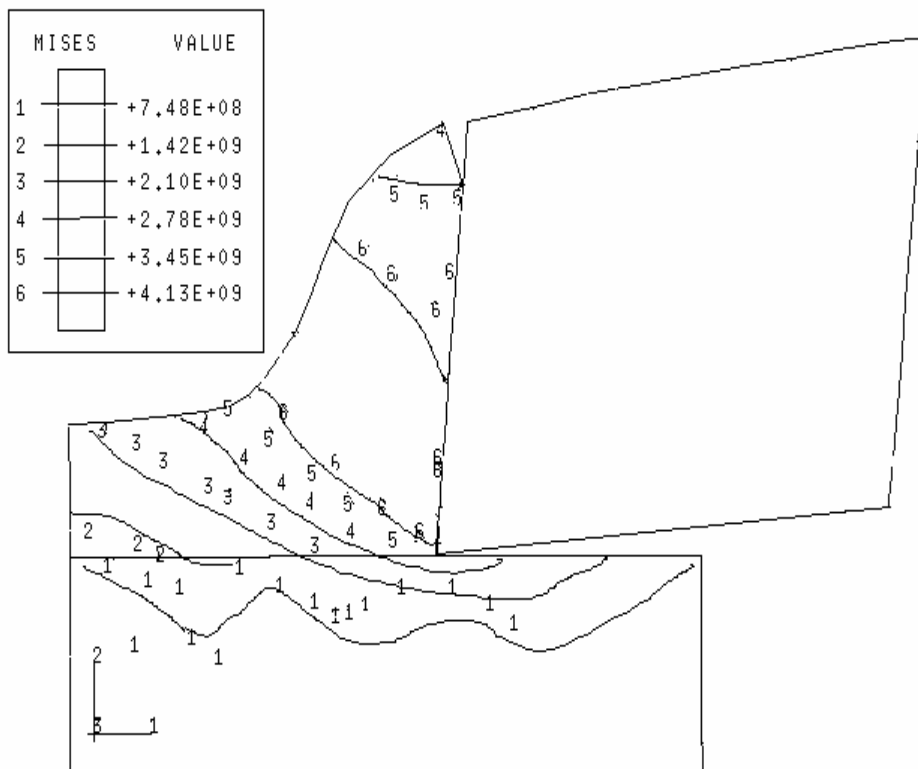


Figure 4. Equivalent Stress

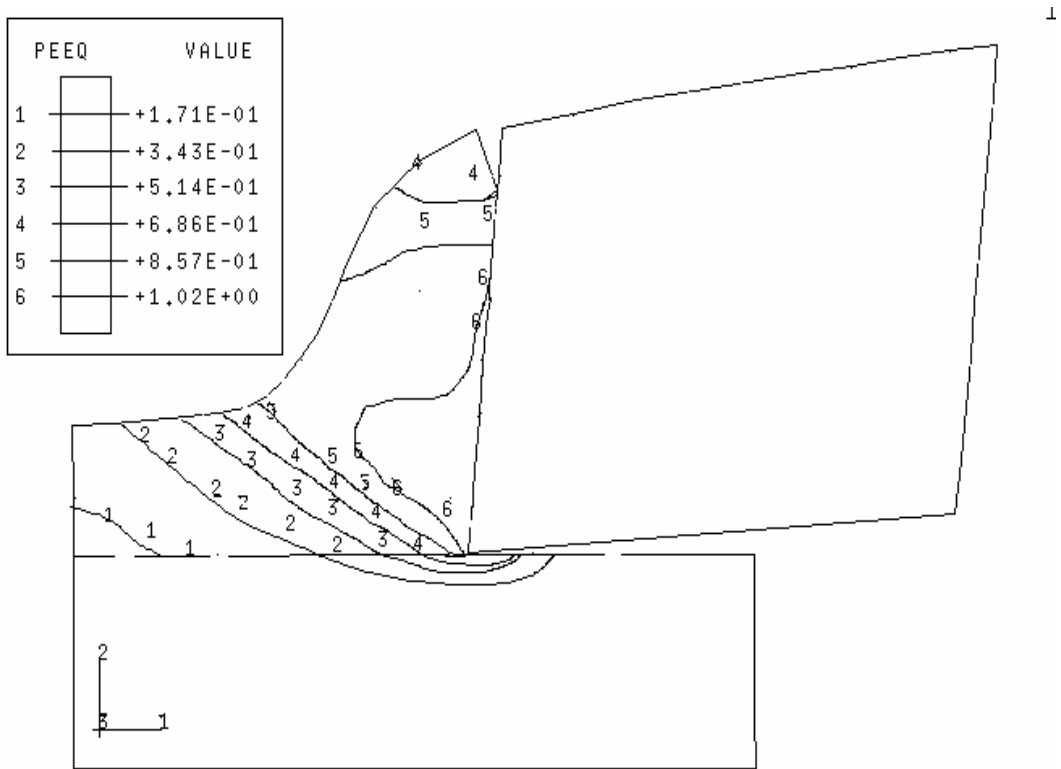


Figure 5. Equivalent plastic strain

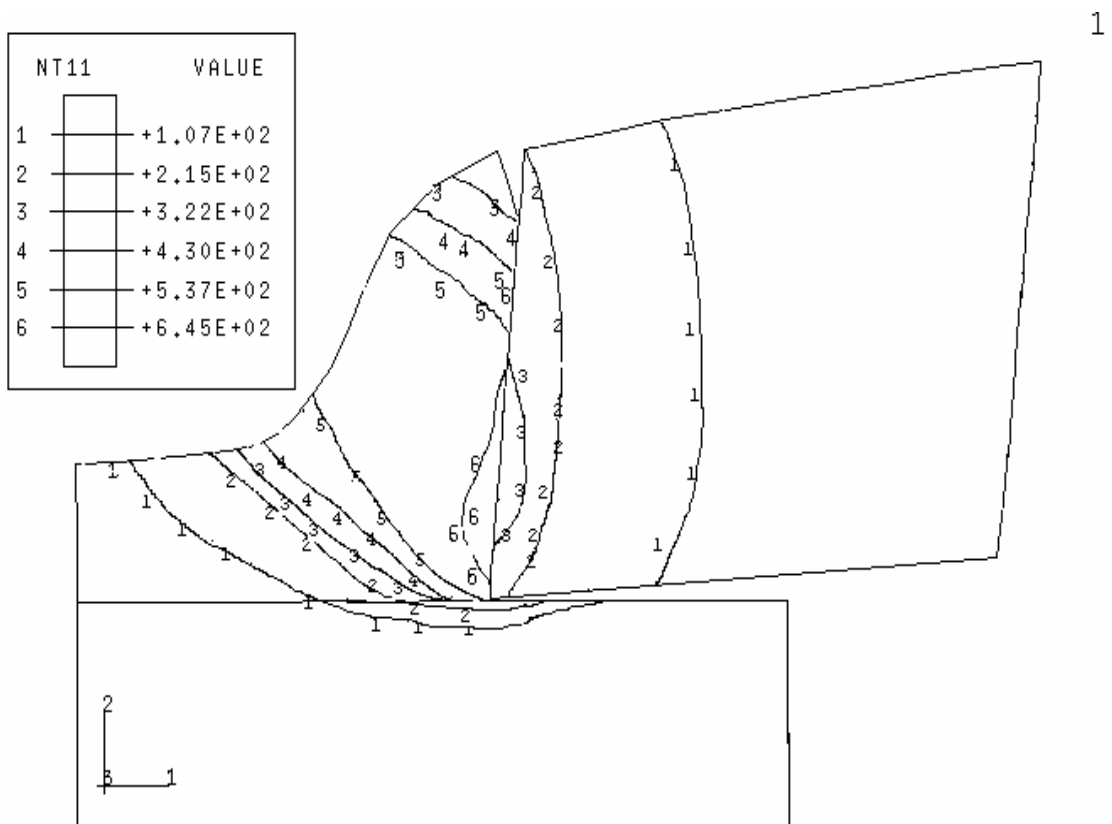


Figure 6. Cutting temperature distribution

4. CONCLUSIONS

Overall observation of the results produced and comparison with some of the available results in the open literature show good agreement. Temperature fields, as well as stress fields behave like, and have values in the range of the ones obtained experimentally or numerically. In the primary zone, top temperatures in the vicinity of 600 degrees were obtained. Only conduction to the tool was allowed and that in a half-and-half proportion. Top temperatures a little bit lower should be expected that the heat transfer in the supposed isolated portions of the model do occur. Allowance for radiation or exchange with lubricants, when they are present, would include small modifications to the model. In what concerns the stresses, again good results are obtained, with discrepancies of the order of 30%. The model presented could be modified also to include anisotropic effects as well as incorporate cinematic hardening conditions. In it the Shirakashi expression for the description of behaviour of metallic materials in a equivalent setting was used (Usui & Shirakashi, 1982). Forces, can also be computed from this model, by integration of the stresses (Madrigal, Batalha & de Aguiar, 2000), and they present the trend and range of measured results.

5. ACKNOWLEDGEMENTS

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